

# Nonlocal theory for fractional kinetic equations

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## A bit of history (of Modena also)

Consider classical kinetic (or Fokker-Planck) equations

$$\underbrace{(\partial_t + v \cdot \nabla_x)}_{\text{transport}} f = \underbrace{\nabla_v \cdot (A \nabla_v f)}_{\text{diffusion}} + \underbrace{B \cdot \nabla_v f + h}_{\text{source}} \quad \underbrace{(t, x, v)}_{\text{times, space, velocity}} \in \mathbb{R} \times \mathbb{R}^n \times \mathbb{R}^n$$

under natural ellipticity conditions on the coefficients and scalar field

$$\begin{cases} 0 < \lambda \mathbb{I}_n \leq A \leq \Lambda \mathbb{I}_n \\ |B| \leq \Lambda \\ h \text{ essentially bounded} \end{cases}$$

In this scenarion, classical first Hölder regularity in the spirit of the works of De Giorgi-Nash & Moser (*DGNM* in short) theroy is completed

- Pascucci & Polidoro, *CCM* (2004):  $L^\infty$ - $L^2$  estimates via Moser's Iteration.
- Golse, Imbert, Mouhot & Vasseur, *Ann. SNS* (2019): Hölder regularity + Harnack inequality
- Guerand & Mouhot, *JEP* (2021): Weak Harnack inequality.
- Anceschi & Rebusci, *JDE* (2022): Weak regularity for kinetic equations with more than one spatial commutator.
- Anceschi et. al, *Preprint* (2024): Poincaré inequality based on local trajectories and weak Harnack

## What's happen is the diffusion is nonlocal??

We investigate local properties of solutions  $f \equiv f(t, x, v)$  to a wide class of integro-differential equations having as toy model

$$(\partial_t + v \cdot \nabla_x) f + (-\Delta_v)^s f = 0, \quad (t, x, v) \in \mathbb{R} \times \mathbb{R}^n \times \mathbb{R}^n$$

Here we denote with  $(-\Delta_v)^s$  the fractional Laplacian acting only the velocity variables

$$(-\Delta_v)^s f(t, x, v) = p.v. \int_{\mathbb{R}^n} \frac{f(t, x, v) - f(t, x, w)}{|v - w|^{n+2s}} dw \quad s \in (0, 1).$$

What are we looking for?

We want to prove classical results in the spirit of the *DGKM* theory for elliptic and parabolic equations.

# Why are we interested in nonlocal kinetic equations?

Consider the Boltzmann equation

$$(\partial_t + v \cdot \nabla_x) f = \mathcal{Q}(f, f).$$

Under special circumstances the nonlinear and nonlocal collision operator  $\mathcal{Q}(\cdot)$  can be “linearized” as follows

$$\mathcal{Q}(f, f) \approx \mathcal{L}_{K_f} f + \langle \text{lower order terms} \rangle$$

where

$$\mathcal{L}_{K_f} f(t, x, v) := p \cdot v \cdot \int_{\mathbb{R}^n} (f(t, x, w) - f(t, x, v)) K_f(t, x, v, w) dw$$

with the measurable kernel depending on the solution  $f$  itself.

## Question:

Is it possible to apply ideas from the area of integro-differential equations in the context of the Boltzmann equation.

There are several difficulties that we must overcome:

- In what way is the kernel  $K_f$  elliptic?
- Is it possible to extend the regularity results for integro-differential parabolic equations to the setting of the Boltzmann equations?

## Conditional regime

There are  $m_o, M_o, E_o, H_o > 0$  such that for all  $(t, x)$  it holds:

$$\left\{ \begin{array}{l} m_o \leq \int_{\mathbb{R}^n} f(t, x, v) dv \leq M_o \\ \int_{\mathbb{R}^n} f(t, x, v) |v|^2 dv \leq E_o \\ \int_{\mathbb{R}^n} f(t, x, v) \log f(t, x, v) dv \leq H_o \end{array} \right. \quad (H)$$

Under the bounds (H) the operator  $\mathcal{L}_{K_f}$  satisfies:

$$(H) \implies \left\{ \begin{array}{l} \text{Coercivity: } \|f\|_{\dot{H}^s}^2 \lesssim - \int_{\mathbb{R}^n} \mathcal{L}_{K_f} f(v) f(v) dv + \|f\|_{L^2}^2 \\ \text{Non-degeneracy assumption when } 0 < s < 1/2 \\ \text{Weak upper bounds } \int_{\mathbb{R}^n \setminus B_r(v)} K_f(v, w') dw \lesssim r^{-2s} \\ \text{Cancellation conditions} \end{array} \right.$$

## Remark

*The ellipticity constants of  $K_f$  depends on the bounds on the quantities in (H) and on  $\|f\|_{L^\infty}$ .*

## Theorem (Imbert & Silvestre, *JEMS* (2020))

There exists universal constants  $R_o \in (0, 1)$ ,  $R_1 > 1$  and  $\zeta \in (0, 1)$  such that if  $f$  is a nonnegative weak supersolutions to

$$(\partial_t + v \cdot \nabla_x) f = \mathcal{L}_{K_f} f + h \quad \text{in } (-1, 0] \times B_{R_1^{1+2s}} \times B_{R_1}$$

then

$$\|f\|_{L^\zeta(Q^-)} \lesssim \inf_{Q^+} f + \|h\|_{L^\infty},$$

with

$$Q^+ := (-R_o^{2s}, 0] \times B_{R_o^{1+2s}} \times B_r$$

$$\text{and } Q^- := (-1, -1 + R_o^{2s}] \times B_{R_o^{1+2s}} \times B_{R_o}$$

- The weak Harnack inequality is enough to derive Hölder regularity for solutions:

$$\|f\|_{C^\alpha((-1/2, 0) \times B_{1/2} \times B_{1/2})} \lesssim \|f\|_{L^\infty((-1, 0] \times B_1 \times \mathbb{R}^n)} + \|h\|_{L^\infty((-1, 0] \times B_1 \times B_1)}$$

- Both the full Harnack inequality and a linear  $L^\infty$ - $L^2$  do not hold for nonlocal equations as the one modeling the Boltzmann equation.

# Nonlocal Fokker-Planck equations

Consider a “better” equation

$$(\partial_t + v \cdot \nabla_x) f = \mathcal{L}_K f + h \quad (1)$$

with the measurable kernel  $K : \mathbb{R}^{1+3n} \rightarrow [0, +\infty]$  satisfying

$$\begin{cases} K(t, x, v, w) = K(t, x, w, v), \\ K(t, x, v, w) \approx |v - w|^{-n-2s}. \end{cases}$$

- [Stokols, SIMA \(2019\)](#): Hölder regularity for essentially bounded solutions
- [Loher, JFA \(2024\)](#): “Not-so-Strong” Harnack for essentially bounded solutions

$$\sup_{Q^-} f \lesssim \left( \inf_{Q^+} f \right)^\beta \quad \beta \in (0, 1)$$

- [Anceschi et al. arXiv \(2024\)](#): weak Harnack inequality via Poincaré inequality.

## Remark

*The classical full Harnack inequality is absent in all of these previous works*

## Motivation:

Develop the nonlocal theory of Kolmogorov equations completing the key missing results (full Harnack inequality and  $L^\infty$ - $L^2$  estimate).

## Remark

*The transport and diffusion defines a geometric structure which leaves them invariant and preserves their homogeneity.*

Endow  $\mathbb{R}^{1+2n}$  with the following Galilean transformation and scalings

$$\begin{cases} z_o \circ z := (t + t_o, x + x_o + tv_o, v + v_o) \\ D_R(z) := (R^{2s}t, R^{1+2s}x, Rv) \quad \forall R > 0 \end{cases}$$

We introduce a family of domains respecting the invariant transformations defined above. Given

$$Q_1 \equiv Q_1(0) := U_1(0, 0) \times B_1(0) = (-1, 0] \times B_1(0) \times B_1(0).$$

define the *slanted* cylinder  $Q_R(z_o)$  by

$$\begin{aligned} Q_R(z_o) &:= \{z_o \circ D_R(z) : z \in Q_1\} \\ &\equiv \{(t, x, v) : t_o - R^{2s} < t \leq t_o, \\ &\quad |x - x_o - (t - t_o)v_o| < R^{1+2s}, |v - v_o| < R\}. \end{aligned}$$



## Functional setting

For  $s \in (0, 1)$  we denote with  $W^{s,2}(\mathcal{A})$  the classical fractional Sobolev space

$$W^{s,2}(\mathcal{A}) := \left\{ f \in L^2(\mathcal{A}) : [f]_{s,2;\mathcal{A}} < +\infty \right\},$$

where

$$[f]_{s,2;\mathcal{A}} := \left( \int_{\mathcal{A}} \int_{\mathcal{A}} \frac{|f(v) - f(w)|^2}{|v - w|^{n+2s}} dv dw \right)^{1/2},$$

equipped with the usual norm

$$\|f\|_{W^{s,2}(\mathcal{A})} := \|f\|_{L^2(\mathcal{A})} + [f]_{s,2;\mathcal{A}}.$$

**Definition (Di Castro, Kuusi & Palatucci, *JFA* (2014) + *Ann. IHP-C* (2016))**

*Let  $f$  be a measurable function on  $(t_1, t_2) \times \Omega_x \times \mathbb{R}^n \subset \mathbb{R}^{1+2n}$ , then the nonlocal tail of  $f$  centered in  $v_o$  and of radius  $r$  is defined by*

$$\text{Tail}(f; v_o, R) := R^{2s} \int_{\mathbb{R}^n \setminus B_R(v_o)} |f(v)| |v_o - v|^{-n-2s} dv.$$

In connections with the nonlocal tail consider the related tail space

$$L_{2s}^1(\mathbb{R}^n) := \left\{ f \in L_{\text{loc}}^1(\mathbb{R}^n) : \int_{\mathbb{R}^n} \frac{|f(v)|}{(1 + |v|)^{n+2s}} dv < \infty \right\}$$

## Weak formulation

Given  $\Omega := (t_1, t_2) \times \Omega_x \times \Omega_v \subset \mathbb{R}^{1+2n}$  denote by

$$\mathcal{W} := \left\{ f \in L^2_{\text{loc}}((t_1, t_2) \times \Omega_x; W^{s,2}_{\text{loc}}(\Omega_v)) \cap L^1_{\text{loc}}((t_1, t_2) \times \Omega_x; L^1_{2s}(\mathbb{R}^n)) \right. \\ \left. : (\partial_t + v \cdot \nabla_x) f \in L^2_{\text{loc}}((t_1, t_2) \times \Omega_x; W^{-s,2}(\mathbb{R}^n)) \right\}.$$

and by  $\mathcal{E}^K(\cdot)$  the nonlocal energy

$$\mathcal{E}^K(f, \phi) := \iint_{\mathbb{R}^n \times \mathbb{R}^n} (f(v) - f(w)) (\phi(v) - \phi(w)) K(v, w) dv dw.$$

### Definition

A function  $f \in \mathcal{W}$  is a weak subsolution ( resp., supersolution) to (1) in  $\Omega$  if

$$\int_{t_1}^{t_2} \int_{\Omega_x} \mathcal{E}^K(f, \phi) dx dt + \int_{\Omega} (\partial_t f + v \cdot \nabla_x f) \phi dz \stackrel{(\geq \text{ resp. })}{\leq} \int_{\Omega} h \phi dz$$

for any nonnegative  $\phi \in L^2_{\text{loc}}((t_1, t_2) \times \Omega_x; W^{s,2}(\mathbb{R}^n))$  such that  $\phi(t, x, \cdot)$  is compactly supported in  $\Omega_v$ . A function  $f \in \mathcal{W}$  is a weak solution if it is both a weak sub- and supersolution.

# A roadmap to get the full Harnack

Consider

$$(-\Delta_v)^s f = 0 \quad \text{in } B_{2R}(0) \subset \mathbb{R}^n. \quad (2)$$

To get the full Harnack inequality for solutions  $f \geq 0$  to (2) it suffices to combine that subsolutions satisfies

$$\sup_{B_{R/2}(0)} f_+ \leq c(\delta) \|f_+\|_{L^2(B_R(0))} + \delta \text{Tail}(f_+; 0, R/2) \quad \forall \delta \in (0, 1],$$

whereas nonnegative supersolutions satisfies

$$\text{Tail}(f_+, 0, R) \leq c \sup_{B_{3R/2}(0)} f.$$

By a covering argument and choosing  $\delta > 0$  small enough it is possible to reabsorb the tail term and to prove the full Harnack inequality after combination with the weak one

$$\int_{B_R(0)} f(v) dv \leq c \inf_{B_{R/2}(0)} f$$

- Di Castro, Kuusi & Palatucci, *JFA* (2014)
- Cozzi, *JFA* (2017) (via De Giorgi classes).

Similarly happens for parabolic-type problems

- Strömqvist *Ann. IHP-C* (2019) (for **global solutions**)
- Kassmann & Weidner, *Anal. PDE* (2024) + *Duke* (2024)

## Counterexample to the strong Harnack inequality

Theorem (Kassmann & Weidner, *Adv. Math.* (2024))

There exist a constant  $c_o > 0$  such that for every  $\varepsilon \in (0, \frac{1}{4})$  there exists a solution  $f_\varepsilon : \mathbb{R}^{2n} \mapsto [0, 1]$  to

$$v \cdot \nabla_x f_\varepsilon + (-\Delta_v)^s f_\varepsilon = 0 \quad \text{for } (x, v) \in B_1(0) \times B_1(0),$$

such that for  $\xi := (\frac{1}{2}e_n, 0) \in \mathbb{R}^{2n}$ , it holds

$$f_\varepsilon(\xi) \leq c_o \varepsilon^{n(1+2s)-2s} f_\varepsilon(0).$$

In particular,  $f_\varepsilon(0)/f_\varepsilon(\xi) \rightarrow \infty$  as  $\varepsilon \rightarrow 0$ .

- It is an effect purely originating from the combination of the nonlocality of the diffusion combined with the anisotropy of the transport.
- It is very surprising when compared to all the previous literature dealing with local kinetic equations
- No obstruction to regularity
- The classical  $L^\infty$ - $L^2$ -estimate fails in general. Moreover, the  $L^\infty$ - $L^2$  estimate remains false if an  $L^p$ -norm of the tail is added on the right-hand side if  $p < \frac{n(1+2s)}{2s}$ .

## Theorem (Aneschi, Palatucci &amp; Pic., preprint (2025))

Let  $\Omega := (t_1, t_2) \times \Omega_x \times \Omega_v \subset \mathbb{R}^{1+2n}$  be a domain and  $s \in (0, 1)$ . Assume that  $f \in \mathcal{W}$  is a weak subsolution to

$$(\partial_t + v \cdot \nabla_x) f = \mathcal{L}_K f + h \quad \text{in } \Omega.$$

Then, there exists  $p^* \equiv p^*(n, s) > 2$  such that if, for some  $p > p^*$ , it holds  $\text{Tail}(f_+; B) \in L^p_{\text{loc}}((t_1, t_2) \times \Omega_x)$ , for all  $B \Subset \Omega_v$ , and  $h \in L^p_{\text{loc}}(\Omega)$ , then, for any  $Q_r(z_0) \Subset \Omega$  and any  $\delta \in (0, 1]$ , it holds

$$\begin{aligned} \sup_{Q_{\frac{r}{2}}(z_0)} f &\leq c \left( \frac{\langle v_0 \rangle}{\delta r^{n+3s}} \right)^\beta \|f_+\|_{L^2(Q_r(z_0))} + \|h\|_{L^p(Q_r(z_0))} \\ &\quad + \delta \|\text{Tail}(f_+; B_{r/2}(v_0))\|_{L^p(U_r(t_0, x_0))}, \end{aligned} \quad (3)$$

where  $\beta \equiv \beta(n, s, p) > 0$  and  $c \equiv c(n, s, \Lambda, p) > 0$ . Moreover,  $\beta, c \nearrow \infty$  as  $p \searrow p^*$ .

- We want to build a De Giorgi-type argument for proving the supremum estimate.
- The procedure is based on a combination of a Sobolev inequality and an energy estimate.

# How do we gain integrability?

## Remark

*Because of the strong degeneracy of the involved equations we can not rely on embedding theorems given by the function space  $\mathcal{W}$  itself.*

- We rely on the fundamental solutions of the fractional Kolmogorov equation to transfer regularity in all the variables

## Proposition

*For any  $\sigma > 0$ , let  $g$  be a weak solution of*

$$\begin{cases} (\partial_t + v \cdot \nabla_x)g + (-\Delta_v)^s g \leq (-\Delta_v)^{s/2} h_1 + h_2 & \text{in } [-\sigma^{2s}, 0] \times \mathbb{R}^{2n}, \\ g(-\sigma^{2s}, x, v) = g_0(x, v) & \text{in } \mathbb{R}^{2n}. \end{cases}$$

*Assume  $h_1, h_2 \in L^2([-\sigma^{2s}, 0] \times \mathbb{R}^{2n})$ . Then*

$$\|g\|_{L^q([-\sigma^{2s}, 0] \times \mathbb{R}^{2n})} \lesssim \|g_0\|_{L^2(\mathbb{R}^{2n})} + \|h_1\|_{L^2([-\sigma^{2s}, 0] \times \mathbb{R}^{2n})} + \|h_2\|_{L^2([-\sigma^{2s}, 0] \times \mathbb{R}^{2n})}$$

*for any  $2 \leq q \leq 2 + \frac{2s}{n(1+s)}$ .*

## Proof of the gain integrability

Assume  $z_o = 0$ . Fix  $0 < \varrho < r < 1$ , and define  $\sigma := \varrho + (r - \varrho)2^{-3}$  and

$$\begin{cases} \psi = \psi(x, v) \in C_c^\infty(B_{((\varrho+\sigma)/2)^{1+2s}} \times B_{(\varrho+\sigma)/2}) \\ \psi \equiv 1 \text{ on } B_{\varrho^{1+2s}} \times B_\varrho \text{ and } 0 \leq \psi \leq 1, \\ |\nabla_v \psi| \leq c/(r - \varrho) \text{ and } |(v + v_o) \cdot \nabla_x \psi| \leq c\langle v_o \rangle / (r - \varrho)^{1+2s} \end{cases}$$

Then, the function  $g := (f - \kappa)_+ \psi$  satisfies

$$\begin{aligned} & (\partial_t + v \cdot \nabla_x)g + (-\Delta_v)^s g \\ & \leq (-\Delta_v)^s g + (f - \kappa)_+ (v \cdot \nabla_x \psi) + \psi \chi_{\{f > \kappa\}} h + \mathcal{L}_K g \\ & \quad + \left( \text{p.v.} \int_{\mathbb{R}^n} \frac{(f - \kappa)_+(w) (\psi(w) - \psi(v))}{|v - w|^{n+2s}} dw \right) \chi_{\{f > \kappa\}} \\ & =: (-\Delta_v)^{s/2} h_1 + h_2. \end{aligned}$$

### Remark

*Observe that the right-hand side involves fractional differentiation with respect to the  $v$ -variable only, and these are the directions where we got some regularity estimates from energy estimates.*

Estimating the right-hand side depending on long and short interactions as well as on the ranges of fractional index  $s \in (0, 1)$ , we get

$$\begin{aligned}
 & \|h_1\|_{L^2([-\sigma^{2s}, 0] \times \mathbb{R}^{2n})}^2 + \|h_2\|_{L^2([-\sigma^{2s}, 0] \times \mathbb{R}^{2n})}^2 \\
 & \leq \frac{c \langle v_o \rangle^2}{(r - \varrho)^{2(n+3s)}} \|(f - \kappa)_+\|_{L^2(Q_r)}^2 + \frac{c}{(r - \varrho)^2} \int_{U_r} [(f - \kappa)_+]_{s, 2, B_\sigma}^2 dx dt \\
 & \quad + \frac{c |Q_r \cap \{f > \kappa\}|^{1 - \frac{2}{p}}}{(r - \varrho)^{2(n+2s)}} \|\text{Tail}((f - \kappa)_+; B_r)\|_{L^p(U_r)}^{\frac{2}{p}} \\
 & \quad + c |Q_r \cap \{f > \kappa\}|^{1 - \frac{2}{p}} \|h\|_{L^p(Q_r)}^{\frac{2}{p}},
 \end{aligned}$$

The proof finishes combining the above ones with classical energy estimates

$$\sup_{t \in [-T, 0]} \|g(t, \cdot, \cdot)\|_{L_{x,v}^2}^2 + \int_{-T}^0 \|g(t, \cdot, \cdot)\|_{L_x^2 H_v^s}^2 dt \lesssim \|g\|_{L^2}^2 + \langle \text{nonlocal tail terms} \rangle$$



## Theorem (Aneschi, Palatucci & Pic., preprint (2025))

Let  $\Omega := (t_1, t_2) \times \Omega_x \times \Omega_v \subset \mathbb{R}^{1+2n}$  be a domain and  $s \in (0, 1)$ . Assume that  $f \in \mathcal{W}$  is a weak subsolution. Then,

$$f \in L_{loc}^q(\Omega) \quad \forall q \in \left[ 2, 2 + \frac{2s}{n(1+s)} \right].$$

Furthermore, for any  $p > 2$  such that  $\text{Tail}(f_+; B) \in L_{loc}^p((t_1, t_2) \times \Omega_x)$ , for any  $B \Subset \Omega_v$  and  $h \in L_{loc}^p(\Omega)$ , and any  $Q_r \Subset \Omega$ , the following estimate does hold, for  $0 < \varrho < r$ ,

$$\begin{aligned} & \| (f - \kappa)_+ \|_{L^q(Q_\varrho)} \\ & \leq \frac{c \langle v_0 \rangle}{(r - \varrho)^{2(n+3s)}} \| (f - \kappa)_+ \|_{L^2(Q_r)} + \frac{c |Q_r \cap \{f > \kappa\}|^{\frac{1}{2} - \frac{1}{p}}}{r - \varrho} \| h \|_{L^p(Q_r)} \\ & \quad + \frac{c |Q_r \cap \{f > \kappa\}|^{\frac{1}{2} - \frac{1}{p}}}{(r - \varrho)^{2(n+2s)}} \| \text{Tail}((f - \kappa)_+; B_r) \|_{L^p(U_r)}, \end{aligned}$$

for any  $\kappa \in \mathbb{R}$ .

## De Giorgi iteration and the supremum estimate

We built a nonlinear recursive argument. For  $j \in \mathbb{N}$  define

$$R_j := \frac{1}{2} \left(1 + \frac{1}{2^j}\right) R \quad \text{and} \quad \kappa_j := \left(1 - \frac{1}{2^j}\right) \kappa_o \quad \text{for } R, \kappa_o > 0$$

and

$$Y_j := \kappa_o^{-2} \int_{Q_{R_j}} (f - \kappa_j)_+^2 \, dz$$

Applying the local gain of integrability we thus arrive at

$$Y_{j+2} \lesssim \delta^{-1} \mathbf{b}^j Y_j^{1+\alpha} \quad \text{with } \mathbf{b} > 1,$$

where

$$\begin{cases} \delta \in (0, 1] & (\text{arbitrary}) \\ \alpha := 1 - \frac{2}{p} - \frac{2}{q} & \text{for } p > 2, \, q \in \left(2, 2 + \frac{2s}{n(1+s)}\right] \end{cases},$$

Choosing

$$\begin{cases} p > 2 + \frac{2n(1+s)}{s} \\ \kappa_o \approx \left( \frac{\langle v_o \rangle}{\delta r^{n+3s}} \right)^{\frac{1}{\alpha}} \|f_+\|_{L^2(Q_r)} + \|h\|_{L^p(Q_r)} + \delta \|\text{Tail}(f_+; B_{r/2})\|_{L^p(U_r)}, \end{cases},$$

we obtain that

$$\alpha > 0 \quad \text{and} \quad Y_o \lesssim \mathbf{b}^{-\frac{2}{\alpha^2}}$$

so that a classical iteration argument yields that  $Y_j \rightarrow 0$  as  $j \rightarrow \infty$ , giving the desired result.

Theorem (Anceschi, Palatucci & Pic., preprint (2025))

Let  $\Omega := (t_1, t_2) \times \Omega_x \times \Omega_v \subset \mathbb{R}^{1+2n}$  be a domain,  $Q_2(0) \Subset \Omega$ , and  $s \in (0, 1)$ . Assume that  $f \in \mathcal{W}$  is a globally nonnegative weak solution to

$$(\partial_t + v \cdot \nabla_x) f = \mathcal{L}_K f + h \quad \text{in } \Omega.$$

Then, there exists  $p^* \equiv p^*(n, s) > 2$  such that if, for some  $p > p^*$ , it holds  $\text{Tail}(f_+; B) \in L^p_{\text{loc}}((t_1, t_2) \times \Omega_x)$ , for all  $B \Subset \Omega_v$ , then there exists  $R_o \in (0, 1)$  depending only on  $n$  and  $s$  such that

$$\sup_{Q^-} f \lesssim \inf_{Q^+} f + \|\text{Tail}(f; 0, R_o/2)\|_{L^p(U_{R_o}(-1+R_o^{2s}, 0))},$$

where

$$Q^+ := (-R_o^{2s}, 0] \times B_{R_o^{1+2s}} \times B_{R_o}$$

$$\text{and } Q^- := (-1, -1 + R_o^{2s}] \times B_{R_o^{1+2s}} \times B_{R_o}$$

# Proof of the strong Harnack inequality

We combine the weak Harnack inequality with the following covering Lemma

## Lemma

*There exist constants  $c_* \equiv c_*(s) \in (0, 1)$  and  $\gamma \equiv \gamma(s) \geq 1$  such that, for any  $1/2 \leq \varrho < r \leq 1$  and any  $z_o \in \mathbb{R}^{1+2n}$ , it holds*

$$Q_{c_*(r-\varrho)\gamma}(z) \subset Q_r(z_o) \quad \forall z \in Q_\varrho(z_o),$$

For  $1/2 \leq \sigma' < \sigma \leq 1$ , by the  $L^\infty$ - $L^2$  estimate we have, for  $\gamma_1, \gamma_2 > 0$  (universal)

$$\begin{aligned} \sup_{Q_{\sigma'R_o}(-1+R_o^{2s}, 0, 0)} f &\leq \frac{c(\delta)\|f\|_{L^\zeta(Q_{R_o}^-)}}{[(\sigma - \sigma')R_o]^{\gamma_1}} + \left(c\delta + \frac{2 - \zeta}{2}\right) \sup_{Q_{\sigma R_o}(-1+R_o^{2s}, 0, 0)} f \\ &\quad + \frac{c(\delta)}{[(\sigma - \sigma')R_o]^{\gamma_2}} \|\text{Tail}(f_+; 0, R_o/2)\|_{L^p(U_{R_o}(-1+R_o^{2s}, 0))} \end{aligned}$$

Choosing  $\delta \in (0, 1)$  sufficiently small, we reabsorb the supremum on the left-hand side and conclude with the weak Harnack inequality.

- Our Harnack formulation does not contradict the counterexample built via the sequence  $\{f_\varepsilon\}$  since

$$\frac{\sup f_\varepsilon}{\inf f_\varepsilon + \|\text{Tail}(f_\varepsilon)\|_{L^p}} < \infty \quad \text{as } \varepsilon \searrow 0.$$

- The lower bound of the exponent  $p$  is given by the highest integrability range achievable via convolution with the fundamental solution.
- Weak solutions are not required to have finite  $p$ -tail. However, in accordance with the Boltzmann case, boundedness on the mass trivially implies finite  $p$ -tail, as, e. g.
  - [Silvestre CMP \(2016\)](#): Theorem 1.1 – 1-2
  - [Imbert, Mouhot & Silvestre, JEP \(2020\)](#): Formula (1.3)
  - [Imbert & Silvestre JEMS \(2020\)](#): Section 1.3
  - [Imbert & Silvestre, JAMS \(2022\)](#): Assumption 1.2

THANK YOU